

C. Models of Oscillatory dynamics

1. Realistic predator-prey model

- In Sec A3, we saw that oscillatory sol'n of the Lotka-Volterra model was destroyed when carrying capacity of the prey was included.
(Small prey pop drives predator to extinction)
- observed osc in predator/prey systems?
 - here: include limited "uptake capacity" by predators
 - alternative: stochastic effects at low pop density

$$\frac{dp}{dt} = r p \left(1 - \frac{p}{P_k}\right) - \gamma p \frac{p}{(1 + p/P_k)}$$
$$\frac{dg}{dt} = +\gamma p \frac{p}{(1 + p/P_k)} - \delta g$$

Sec A2: $\gamma p = \text{const}$

Monod form for
"uptake" of prey
by predator

Compared to problems we have analyzed:

- the damped predator-prey system of Sec A3 is obtained by taking $P_k \rightarrow \infty$;
- the predation problem (Sec A2) is obtained by setting $\gamma p = \text{constant}$.

* Make dimensionless (Same notation as in Sec A3)

$$u = p/\tilde{p} \quad v = \frac{rp}{\tilde{p}} \quad , \quad \frac{p_k}{\tilde{p}} = \kappa$$

$$\tilde{\tau} = r \cdot t, \quad \frac{s}{r\tilde{p}} = \gamma$$

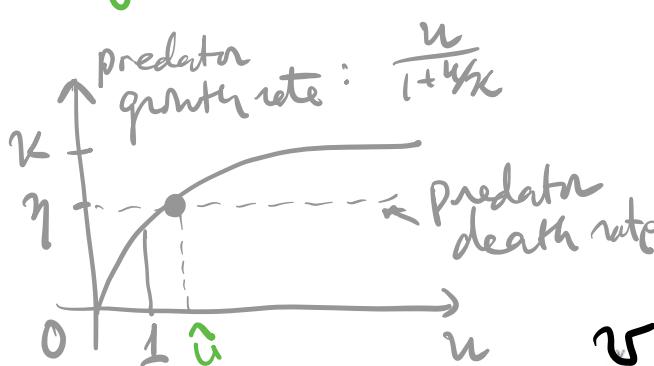
max predator growth rate
(When prey at carrying capacity)

$$\begin{cases} \frac{du}{d\tilde{\tau}} = u(1-u) - \frac{uv}{1+u/\kappa} = f(u, v) \\ \frac{v}{s} \frac{dv}{d\tilde{\tau}} = \frac{v}{\gamma} \left(\frac{u}{1+u/\kappa} - \eta \right) = g(u, v) \end{cases}$$

↑ time scale doesn't affect phase boundary
but affects eigenvalue.

nullclines: $f(u, v) = 0$; $u=0; v=(1-u) \cdot (1+u/\kappa)$

$g(u, v) = 0$ $v=0, u=\eta(1+u/\kappa)$



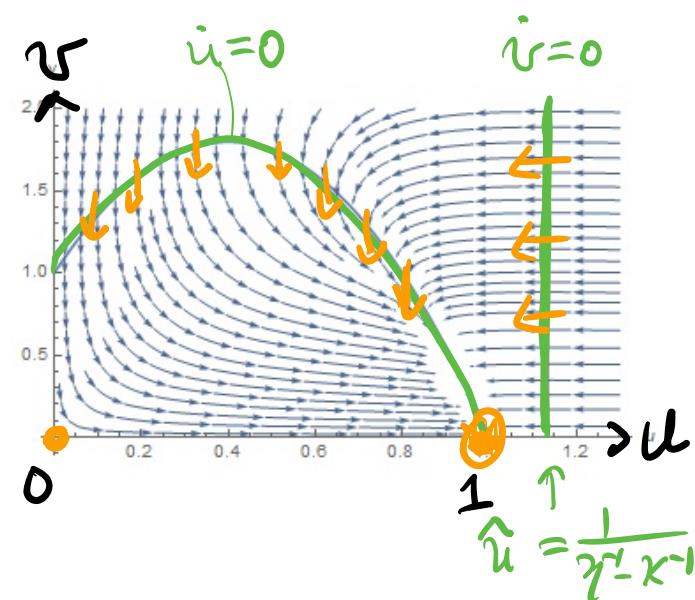
$$\hat{u} = \frac{1}{\gamma^{-1} - \kappa^{-1}}$$

predator extinct
if $\gamma > \kappa$

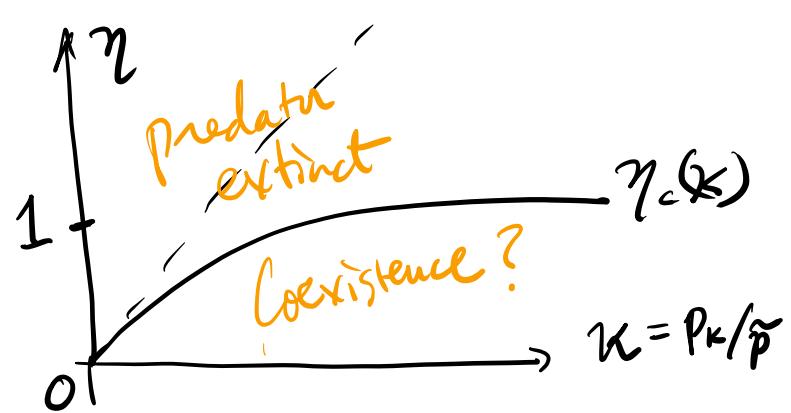
If $\hat{u} > 1$ ($\gamma^{-1} - \kappa^{-1} < 1$, or $\gamma > \frac{\kappa}{\kappa+1}$)

then $u^* = 1, v^* = 0$ is only nontrivial fp.

→ predator extinct,
prey at carrying capacity



* Overview of phase diagram

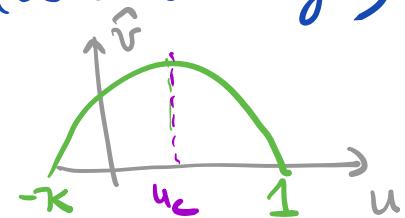


Next, the regime $\gamma < \gamma_c(\kappa) = \frac{\kappa}{1+\kappa}$ where $u^* = \hat{u} = \frac{1}{\gamma^{-1} - \kappa^{-1}} < 1$

\Rightarrow 3 cases depending on the shape of the isocline
 ↗ (controlled by κ)

$$\hat{v}(u) = (1-u) \cdot (1 + \kappa u)$$

$$\text{max: } \left. \frac{d\hat{v}}{du} \right|_{u_c} = 0 \rightarrow u_c = \frac{1-\kappa}{2}$$



Case ①:

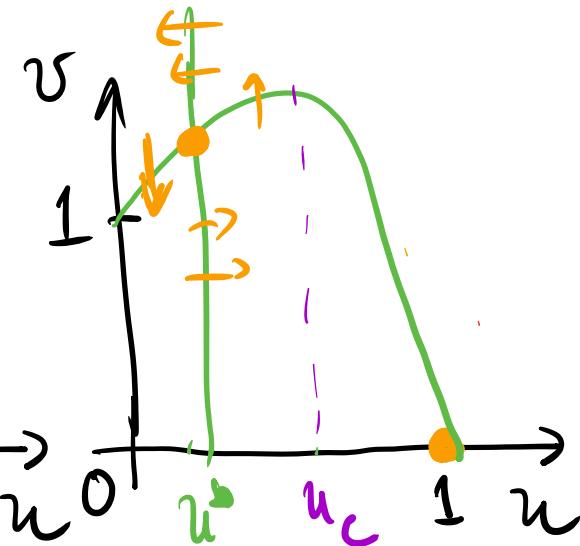
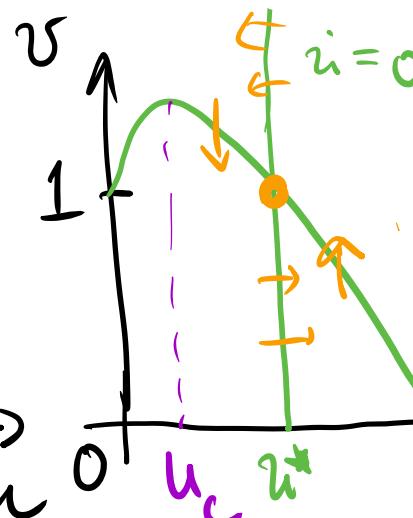
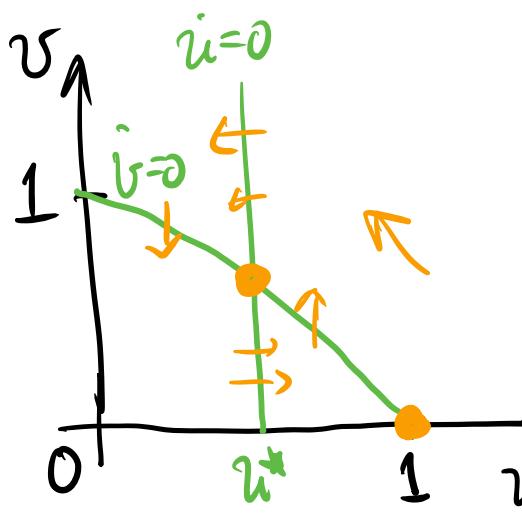
$$\kappa > 1 \rightarrow u_c < 0$$

Case ②:

$$\kappa < 1 \rightarrow u_c > 0 \quad (u^* > u_c)$$

Case ③:

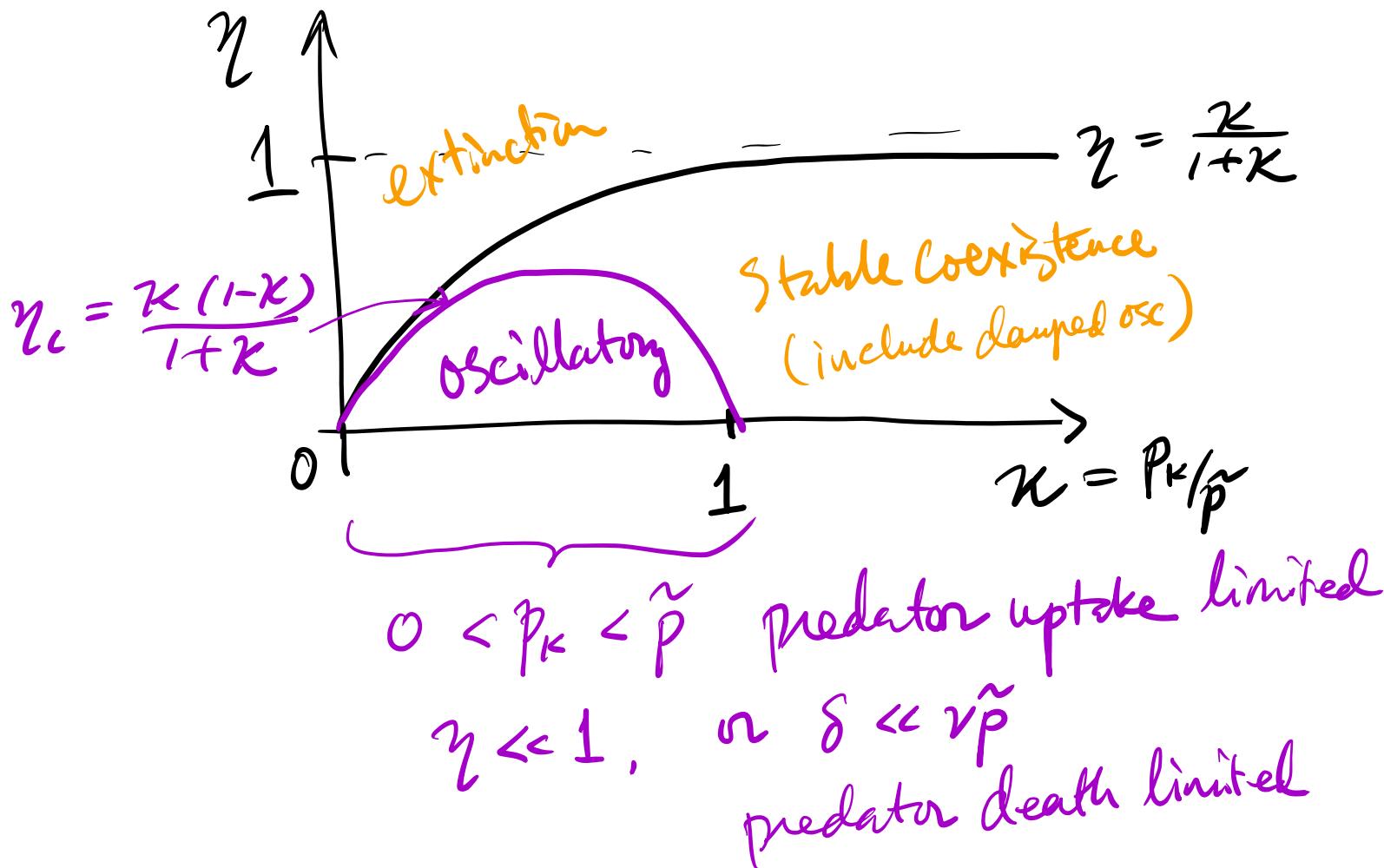
$$\kappa < 1 \rightarrow u_c > 0 \quad (u^* < u_c)$$



$$\rightarrow u^* = u_c \text{ occurs at } \frac{1}{\gamma_c^{-1} - \kappa^{-1}} = \frac{1-\kappa}{2}$$

$$\gamma_c^{-1} = \frac{2}{1-\kappa} + \kappa^{-1} = \frac{1+\kappa^{-1}}{1-\kappa} \quad \text{or} \quad \gamma_c = \frac{\kappa(1-\kappa)}{1+\kappa}$$

will show below that case ③ \rightarrow Stable limit cycle



* Algebraic analysis:

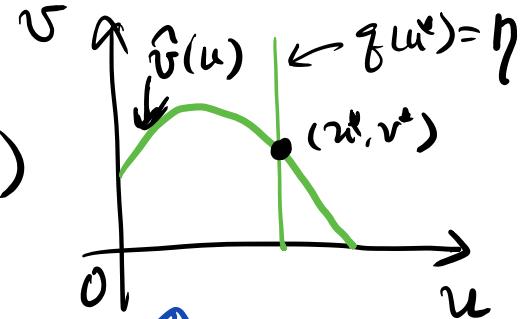
work out the community matrix at fixed pt (w^*, v^*)
Let $u = w^* + x, v = v^* + y$

$$\begin{aligned} \frac{du}{dt} &= f(u, v) \\ \frac{dv}{dt} &= \frac{\delta}{r} g(u, v) \end{aligned} \quad \xrightarrow{\text{linearize}} \quad \begin{pmatrix} \frac{dx}{dt} \\ \frac{dy}{dt} \end{pmatrix} = \underbrace{\begin{pmatrix} \frac{\partial f}{\partial u} & \frac{\partial f}{\partial v} \\ \frac{\delta}{r} \frac{\partial g}{\partial u} & \frac{\delta}{r} \frac{\partial g}{\partial v} \end{pmatrix}}_{M} \cdot \begin{pmatrix} x \\ y \end{pmatrix}$$

$$\left. \begin{array}{l} f(u, v) = p(u) - v q(u) \\ g(u, v) = \left(\frac{1}{2}q(u) - 1\right) \cdot v \end{array} \right| \quad \left. \begin{array}{l} p(u) = u \cdot (1-u) \\ q(u) = u/(1+u/\kappa) \end{array} \right.$$

For the nontrivial fixed pt ($u^* > 0, v^* > 0$)

$$\left. \begin{array}{l} g=0 \rightarrow q(u^*) = \gamma \\ f=0 \rightarrow \hat{v}(u) = \frac{p(u)}{q(u)} = (1-u)(1+\frac{u}{\kappa}) \\ v^* = \frac{p(u^*)}{q(u^*)} = \frac{1}{2}p(u^*) \end{array} \right.$$



Evaluate derivatives at fixed pt:

$$\begin{aligned} \frac{\partial f}{\partial u} &= p' - v^* q' = p' - \frac{p(u^*)}{q(u^*)} q' \\ &= q \cdot \left. \frac{\frac{\partial p}{\partial u} - p \frac{\partial q}{\partial u}}{q^2} \right|_{u^*} = q(u^*) \cdot \left. \frac{d\hat{v}}{du} \right|_{u^*} = \gamma \cdot \left. \frac{d\hat{v}}{du} \right|_{u^*} \end{aligned}$$

$$\frac{\partial f}{\partial v} = -q(u^*) = -\gamma$$

$$\frac{\partial g}{\partial u} = \frac{1}{2} v^* q'$$

$$\frac{\partial g}{\partial v} = \frac{1}{2} q(u^*) - 1 = 0$$

$$M = \begin{pmatrix} \gamma \frac{d\hat{v}}{du}|_{u^*} & -\gamma \\ \frac{1}{2} v^* q' & 0 \end{pmatrix}$$

$$\det(M - \lambda I) = 0 \rightarrow$$

$$\boxed{\lambda^2 - \gamma \frac{d\hat{v}}{du}|_{u^*} \lambda + \frac{1}{4} v^* q' = 0}$$

$$\lambda = \frac{\gamma}{2} \frac{d\hat{v}^*}{du} \pm \sqrt{\left(\frac{\gamma}{2} \frac{d\hat{v}^*}{du} \right)^2 - \Delta} \quad \Delta$$

$$q(u) = \frac{u}{1+u/\kappa}; q' = \frac{1+u/\kappa - u/\kappa}{(1+u/\kappa)^2} > 0 \Rightarrow \Delta < \left(\frac{\gamma}{2} \frac{d\hat{v}^*}{du} \right)^2$$

$$\hat{v}(u) = \frac{p(u)}{q(u)} = (1-u)(1+u/\kappa)$$

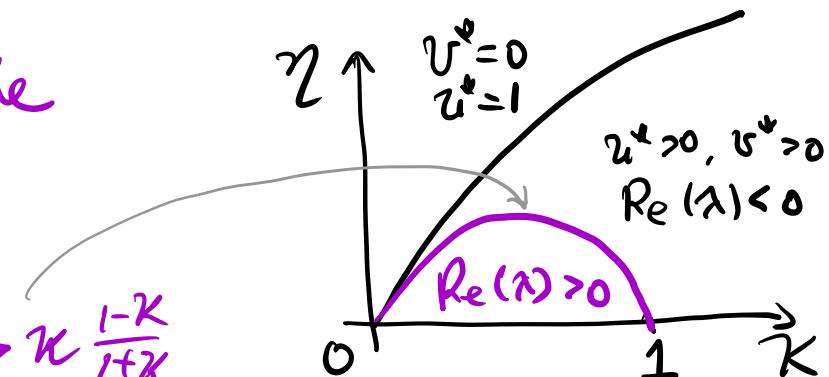
$$\frac{d\hat{v}^*}{du} = \frac{2}{\kappa} \left(\underbrace{\frac{1-\kappa}{2} - u^*}_{u_c} \right) = \frac{2}{\kappa} (u_c - u^*)$$

- If $\frac{d\hat{v}^*}{du} < 0$ (for $u^* > u_c$), $\lambda = - \left| \frac{\gamma}{2} \frac{d\hat{v}^*}{du} \right| \pm \sqrt{\Delta}$, $\operatorname{Re} \lambda < 0$

$\rightarrow (u^*, v^*)$ is stable

Condition: $u^* > u_c$,

$$\frac{1}{2 - \kappa^{-1}} > \frac{1 - \kappa}{2} \rightarrow \gamma > \kappa \frac{1 - \kappa}{1 + \kappa}$$



\rightarrow further calculate Δ to find regimes for
Stable coexistence ($\Delta > 0$) or damped osc ($\Delta < 0$)

- if $\frac{d\hat{v}^*}{du} > 0$ ($u^* < u_c$), $\lambda = \left| \frac{\gamma}{2} \frac{d\hat{v}^*}{du} \right| \pm \sqrt{\Delta} > 0$

\rightarrow to show the limit cycle, need to show $\Delta < 0$

* Calculate the determinant $\Delta = \left(\frac{\gamma}{2} \frac{du^*}{du}\right)^2 - \frac{\delta_r}{r} v^* q'$

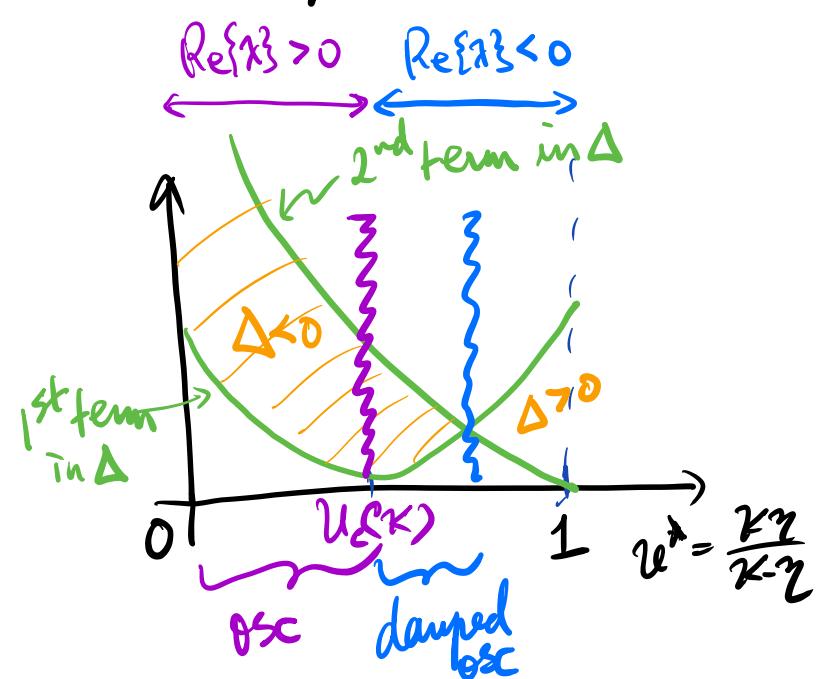
$$\frac{du^*}{du} = \frac{2}{\kappa} (u_c(u) - u^*);$$

$$u^* = \frac{1}{\gamma - \kappa^{-1}} \quad q' = \frac{1}{(1 + u^*/\kappa)^2} \quad \rightarrow q' = \left(\frac{\gamma}{u^*}\right)^2$$

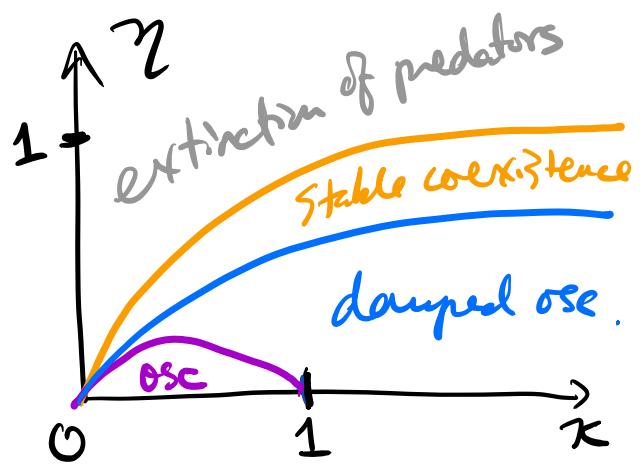
$$v^* = \frac{P(u)}{q(u^*)} = \frac{u^*(1-u^*)}{\gamma}; \quad \text{2nd term in } \Delta = \frac{\delta_r \gamma}{r} \frac{(1-u^*)}{u^*}$$

$$\Rightarrow \Delta = \gamma^2 \left[\left(\frac{u_c(\kappa) - u^*}{\kappa} \right)^2 - \frac{\delta_r \gamma}{r} \frac{1-u^*}{u^*} \right]$$

- To see how Δ depends on u^* , plot each term in [] for fixed κ .



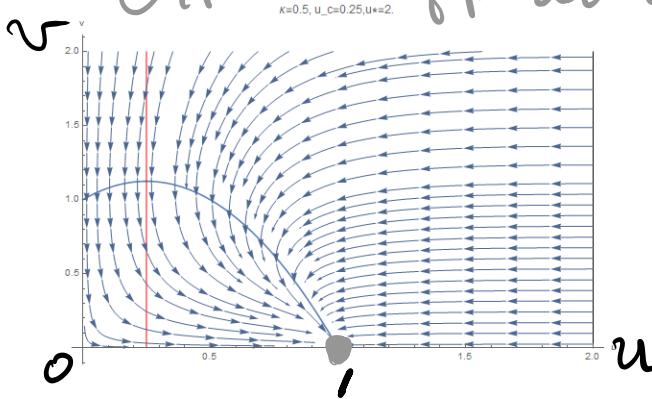
Final phase diagram



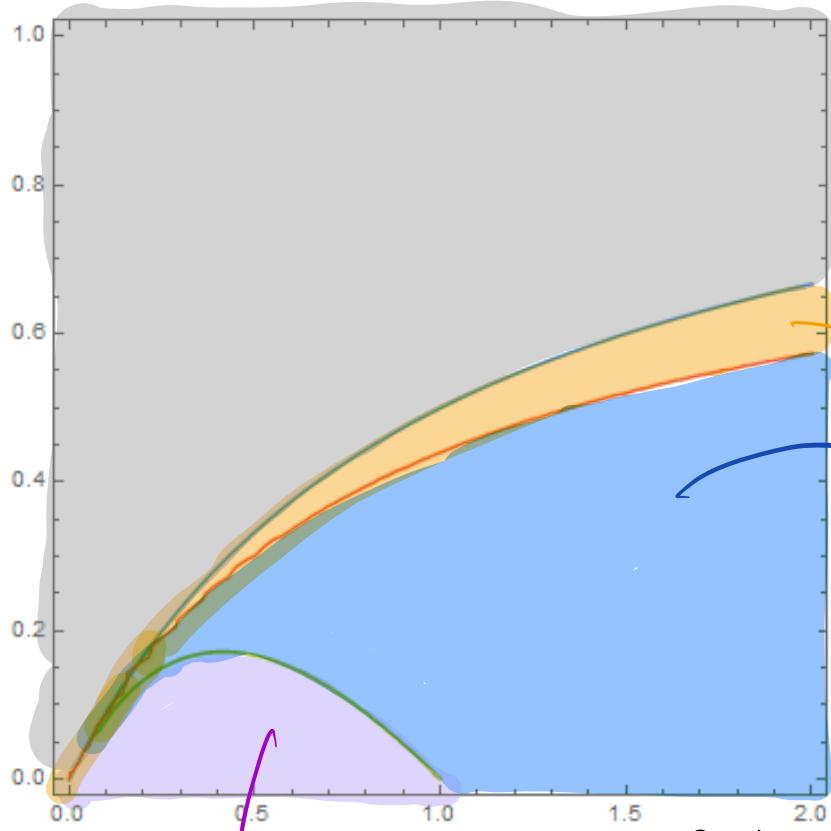
* finally, for large u, v , $\dot{u} = -u^2 - \# u; \dot{v} = -\# v$

Poincaré-Bendixson Theorem: Stable limit cycle for $\Delta < 0$ and $\text{Re}(\lambda) > 0$.

extinction of predator

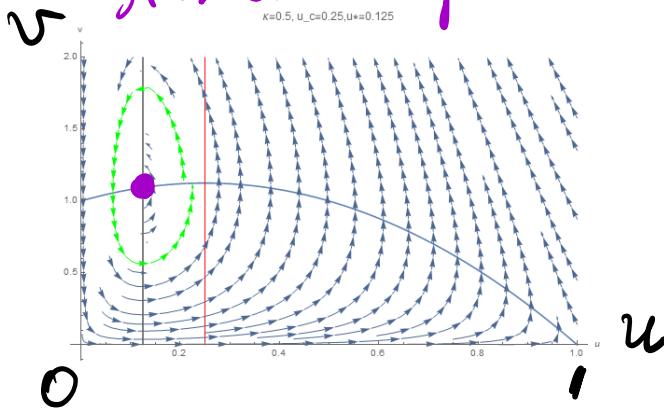


$$\gamma = \delta/\nu^2 = \kappa$$

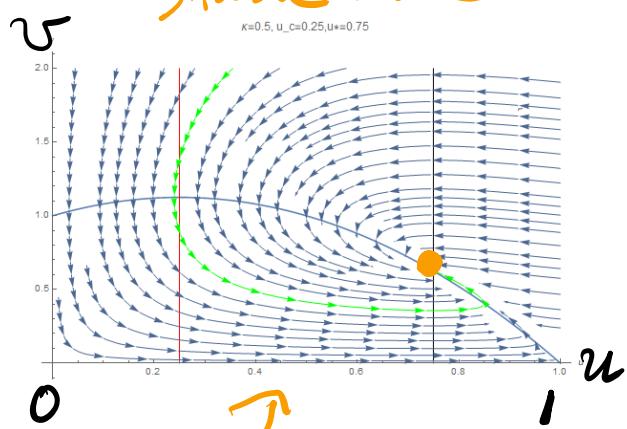


$$\kappa = P_k / \tilde{P}$$

Stable limit cycle



Stable node



clamped oscillation

