

# Tutorial TA session

PHYS 239 – Spatiotemporal Dynamics in Biological Systems

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January 10th, 2022



UC San Diego



- Taylor expansion

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- First-order linear ODEs and coupled linear ODEs

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**Important!**

Please *DO* interrupt me at *any time* if you have questions!





## Basic idea

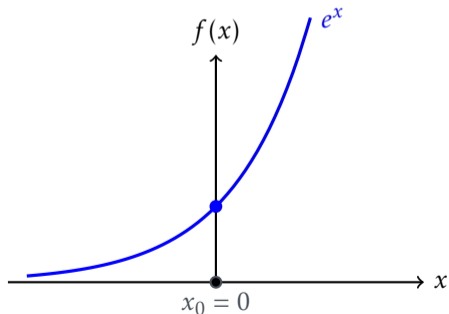
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Why? Because polynomials are simple! This way, we can express complicated functions with simpler terms (at least locally).

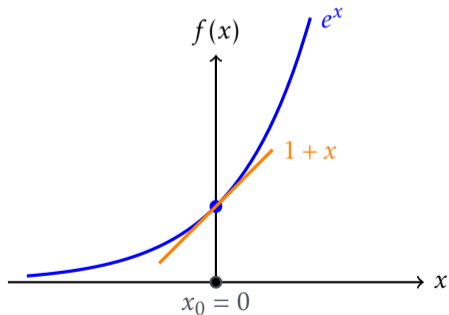
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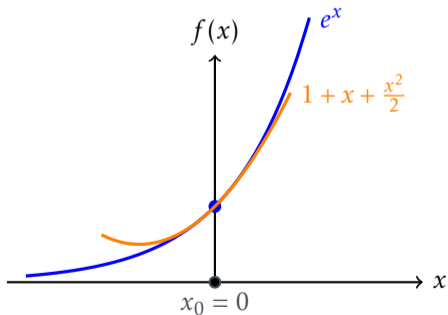
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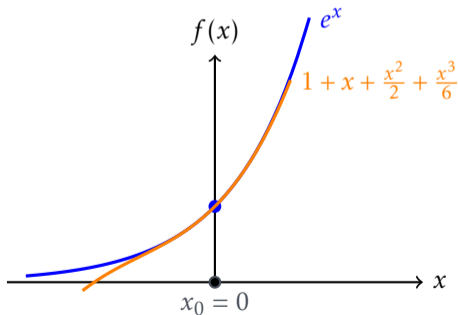
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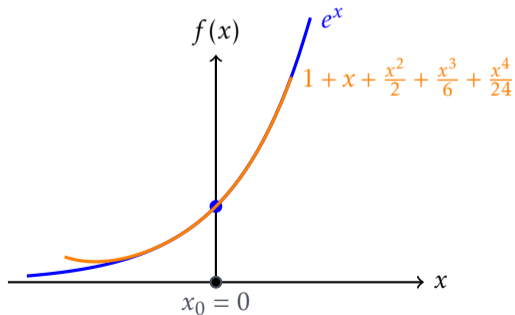
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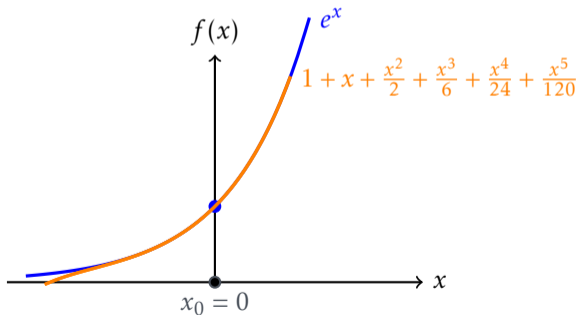
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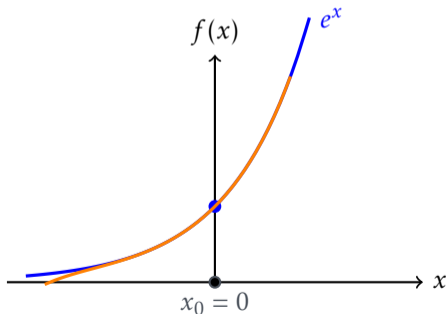
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## Taylor's theorem

If  $f(x)$  is continuous (i.e., it doesn't have any "jumps") and differentiable (i.e., it doesn't have "cusps" or "spikes") in  $x_0$ , then around  $x_0$  we can approximate:

$$f(x) \approx f(x_0) + f'(x_0)(x - x_0) + \frac{1}{2}f''(x_0)(x - x_0)^2 + \frac{1}{6}f'''(x_0)(x - x_0)^3 + \dots \quad (1)$$

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This last point is particularly important, because we can say a lot about linear ODE but almost *nothing* about **non**-linear ODEs in general. Thanks to this theorem, however, we can understand the behavior of nonlinear ODEs around specific points.





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This is just a trick and not the mathematically rigorous way to show that the solution of eq 2 is an exponential. But it works, so we use it.



A system of couple ODEs are two or more ODEs where each variable depends on the other ones.  
For example:

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The trick, in this case, is to use the matrix notation to rewrite the equation in terms of one two-dimensional variables:

$$\Rightarrow \frac{d}{dt} \begin{pmatrix} x \\ y \end{pmatrix} = \underbrace{\begin{pmatrix} a_1 & b_1 \\ a_2 & b_2 \end{pmatrix}}_{:=A} \begin{pmatrix} x \\ y \end{pmatrix} = A \begin{pmatrix} x \\ y \end{pmatrix} \quad (4)$$

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We can now use the same trick shown before to write the solution of eq (3):

$$\begin{pmatrix} x(t) \\ y(t) \end{pmatrix} = \mathbf{exp}(At) \begin{pmatrix} x(0) \\ y(0) \end{pmatrix} \quad (5)$$

even though it is not clear what is the meaning of the exponential of a matrix. We are also assuming  $t_0 = 0$  for the sake of simplicity.

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For our purposes we actually don't need to know what the exponential of a matrix is. A mathematical theorem, in fact, ensures us that we can write the solution of our coupled ODEs as:

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## Therefore

In order to solve a system of coupled linear ODEs we just need to know the so-called *spectral properties* (i.e., eigenvalues and eigenvectors) of the matrix  $A$

As an example, let's solve this system:

$$\frac{dx}{dt} = 3x - 4y \quad \frac{dy}{dt} = 4x - 7y \quad (8)$$

with initial conditions  $x(0) = y(0) = 1$ .

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where  $\mathbb{I} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$  is the identity matrix.

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Therefore, we need to solve:

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$$\det(A - \lambda I) = 0 \quad (11)$$

Therefore, we need to solve:

$$\det\left[\begin{pmatrix} 3 & -4 \\ 4 & -7 \end{pmatrix} - \lambda \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}\right] = \det\begin{pmatrix} 3 - \lambda & -4 \\ 4 & -7 - \lambda \end{pmatrix} = 0 \quad (12)$$

Since by definition:

$$\det\begin{pmatrix} a & b \\ c & d \end{pmatrix} = ad - bc \quad (13)$$

We have:

$$\det\begin{pmatrix} 3 - \lambda & -4 \\ 4 & -7 - \lambda \end{pmatrix} = (3 - \lambda)(-7 - \lambda) - 4(-4) = 0 \Rightarrow -21 - 3\lambda + 7\lambda + \lambda^2 + 16 = 0 \Rightarrow \\ \Rightarrow \lambda^2 + 4\lambda - 5 = 0 \quad (14)$$

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We can now use the eigenvalues to find the eigenvectors.

By definition the eigenvectors  $\vec{u}_1$  and  $\vec{u}_2$  satisfy  $A\vec{u}_1 = \lambda_1\vec{u}_1$  and  $A\vec{u}_2 = \lambda_2\vec{u}_2$ , respectively.

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Therefore, the solution of our system of coupled ODEs is:

$$x(t) = \frac{2}{3} e^t + \frac{1}{3} e^{-5t} \quad y(t) = \frac{1}{3} e^t + \frac{2}{3} e^{-5t} \quad (24)$$





We would like to solve non-linear ODEs (and systems of coupled ODEs), i.e.:

$$\dot{x} = f(x) \tag{25}$$

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Contrarily to linear ODEs, there is no general theorem that allows us to solve non-linear ODEs in general. They almost always *cannot* be solved analytically.

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## Can we say anything about these systems without solving them analytically?

We can use Taylor expansion to approximate the non-linear function  $f(x)$  with a linear one around a point of interest  $x_0$ , and solve the linearized ODEs.

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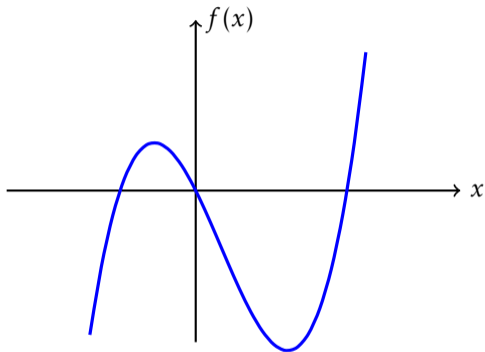
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We can get a sense of the stability of equilibria without even trying to solve a non-linear system by drawing *stream plots* (or *flow plots*).



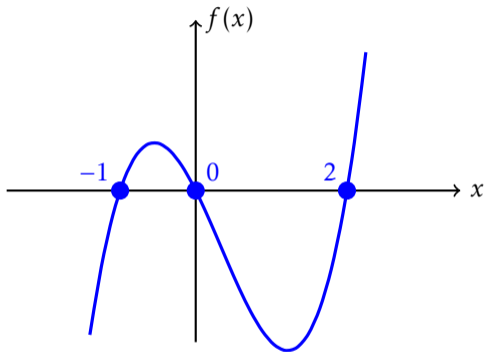
Let's see how to do this in a particular case:  $\dot{x} = x^3 - x^2 - 2x$

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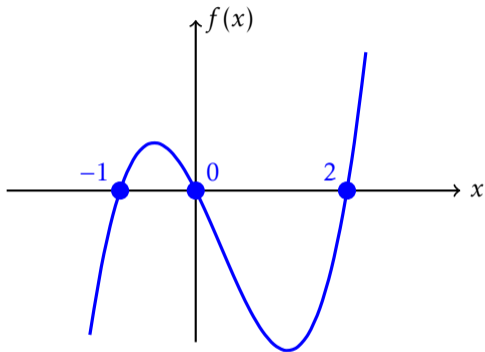


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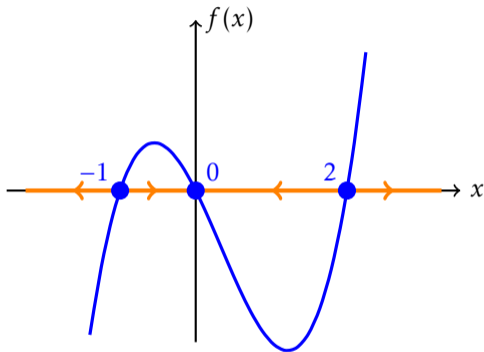
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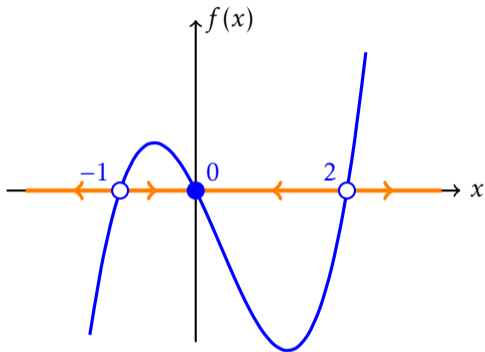
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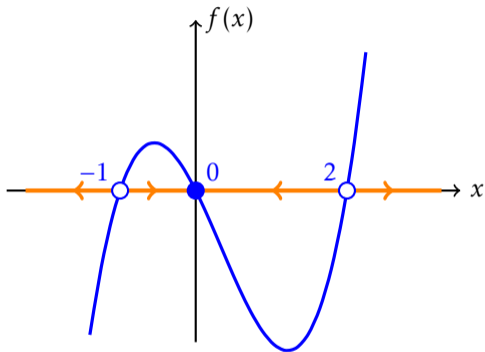
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## Important

In one dimension, if there is more than one equilibrium their stability always “alternates”: after a stable equilibrium we must find an unstable one, and viceversa.



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If  $f(\vec{x}_0) = 0$  (i.e.,  $\vec{x}_0$  is an equilibrium) we can use Taylor's theorem to approximate the system around  $\vec{x}_0$  as:

$$\dot{x} = \underbrace{f(\vec{x}_0)}_{=0} + J(\vec{x}_0)(\vec{x} - \vec{x}_0) = J(\vec{x}_0)(\vec{x} - \vec{x}_0) \quad (30)$$

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$$\vec{x}(t) = \vec{x}_0 + \sum_{i=1}^n c_i e^{\lambda_i t} \vec{u}_i \quad (33)$$

where  $\lambda_i$  are the eigenvalues of  $J(\vec{x}_0)$ ,  $\vec{u}_i$  are its eigenvectors, and  $c_i$  are constants (to be determined from the initial conditions).

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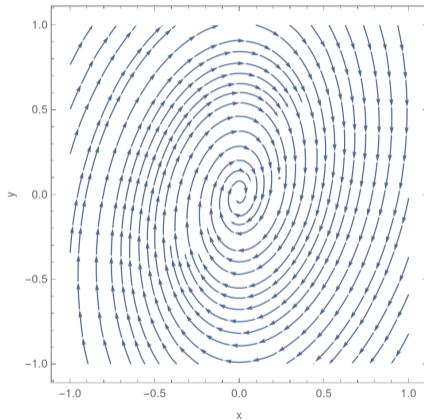
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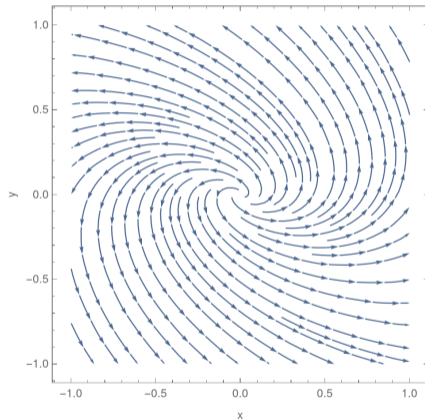
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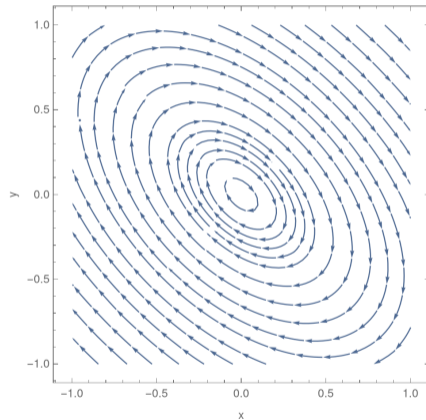
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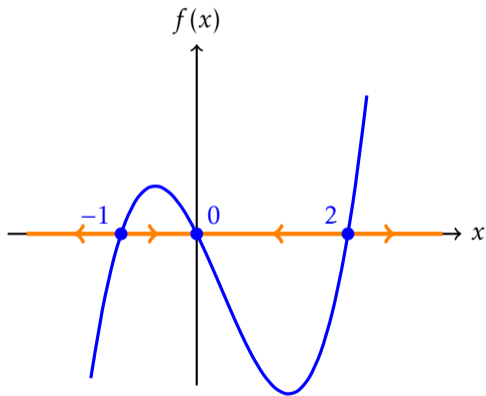
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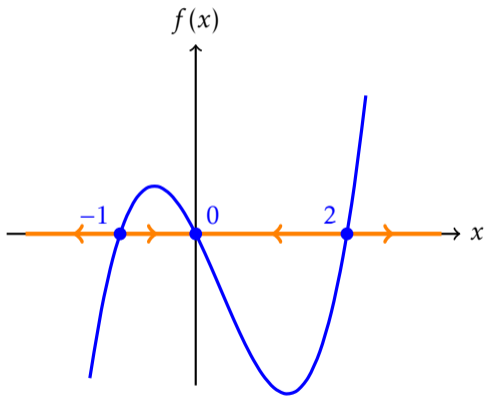
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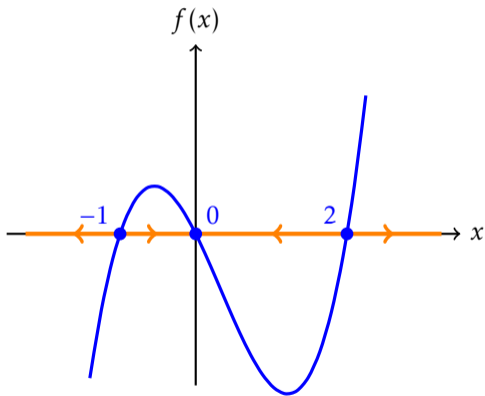
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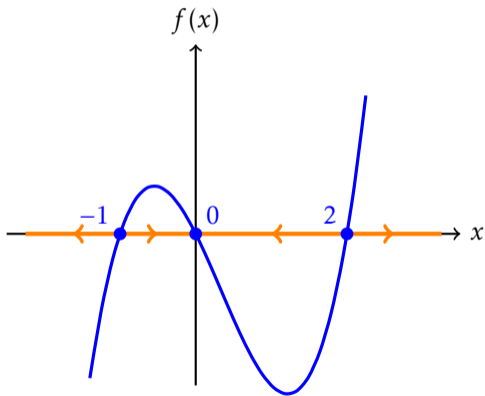
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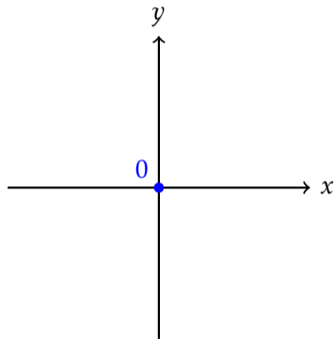
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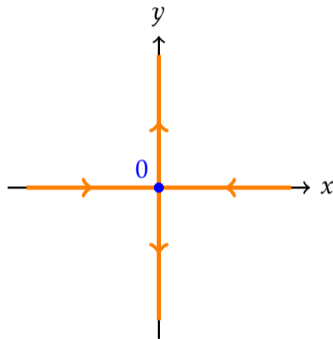
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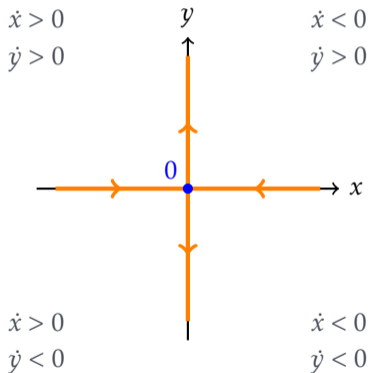
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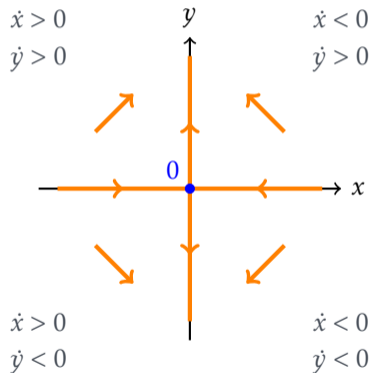
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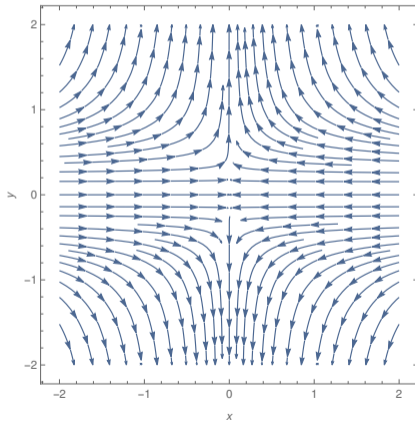
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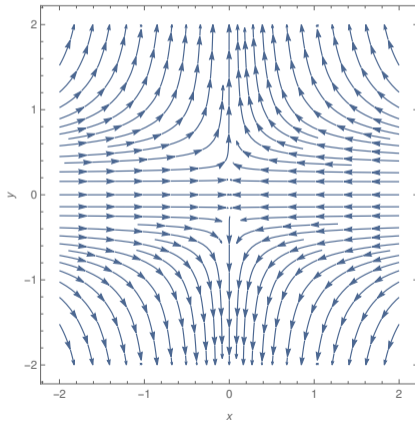
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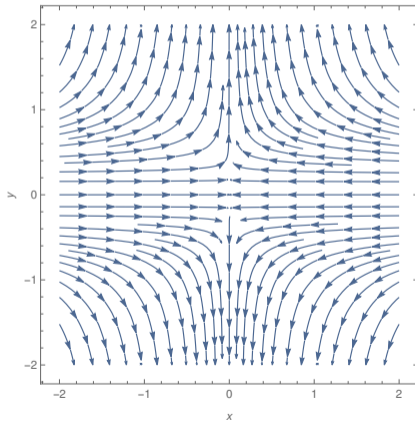
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*Excercise: what happens if  $k < 0$ ?*







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In general, as the parameters of a dynamical system are changed, equilibria can be created or destroyed, or their stability can change.

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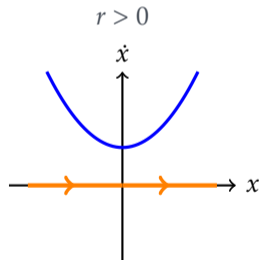
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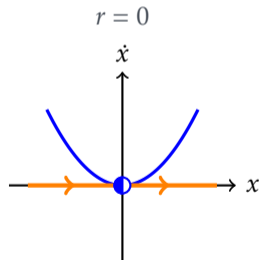
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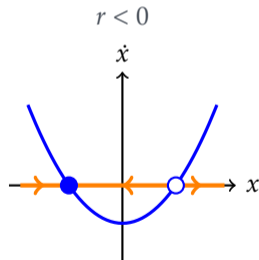
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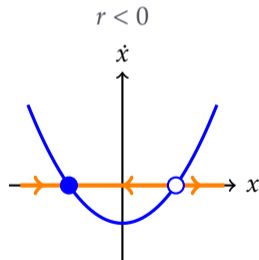
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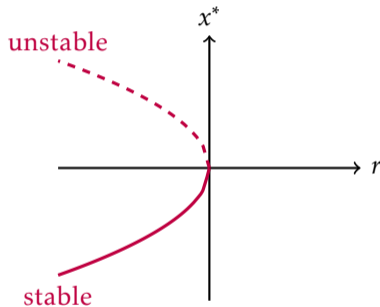
The bifurcation point in this system is  $r = 0$ .



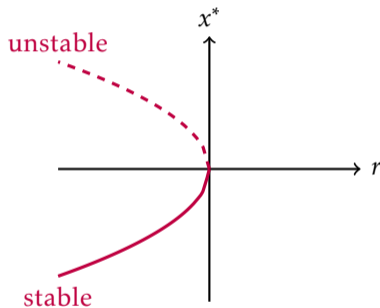


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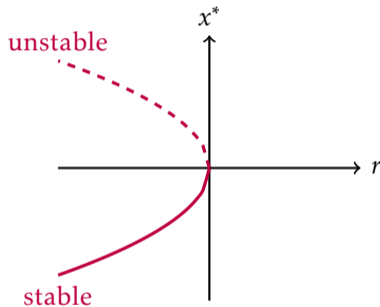


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This type of bifurcation is also known as *turning point bifurcation*, because the bifurcation point  $r = 0$  can also be called *turning point*



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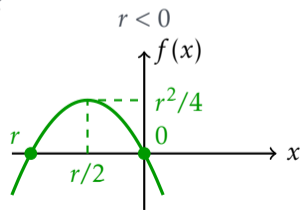
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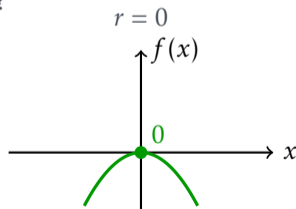
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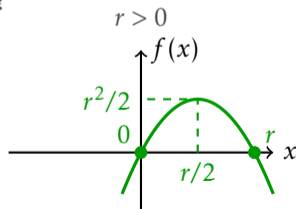
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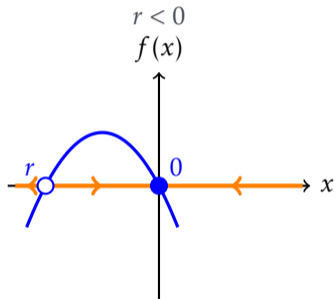
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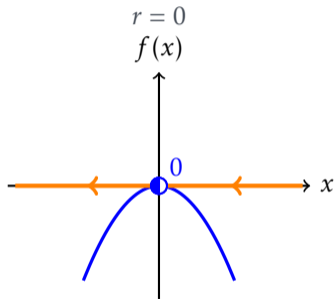
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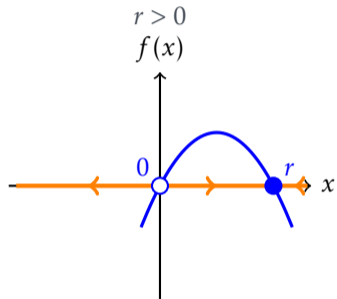


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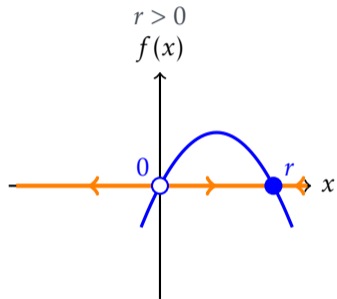


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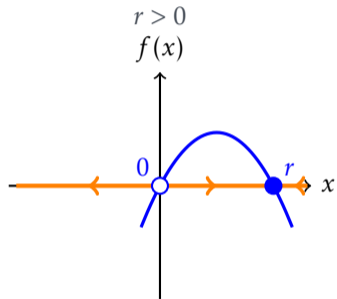
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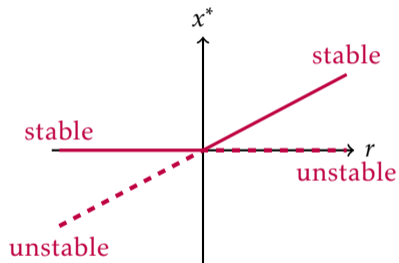


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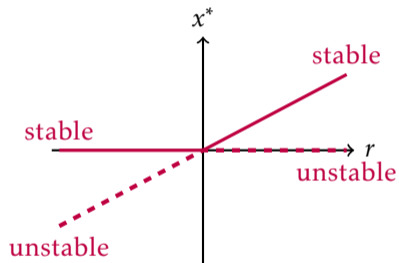
One equilibrium always remains in  $x^* = 0$ , but as  $r$  is changed another equilibrium ( $x^* = r$ ) “crosses” over it and “exchanges stability” with it.

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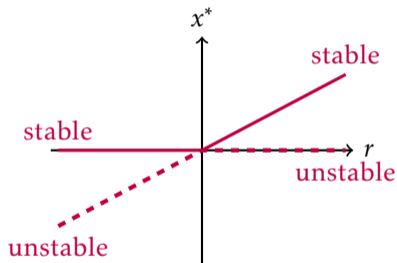


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Contrarily to the saddle-point bifurcation no equilibrium is created or destroyed in this case, but the bifurcation leads to the stability being “exchanged” between equilibria.

# Bifurcation theory

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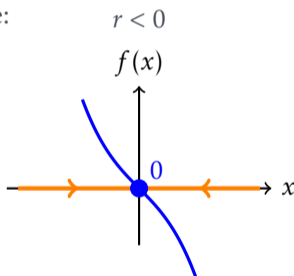
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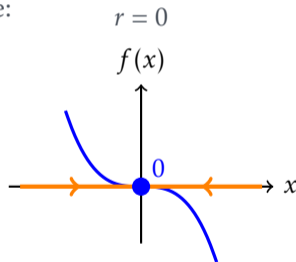


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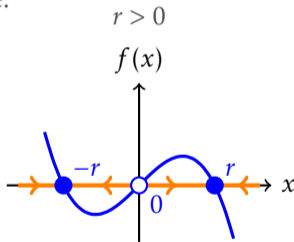


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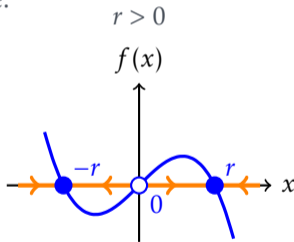


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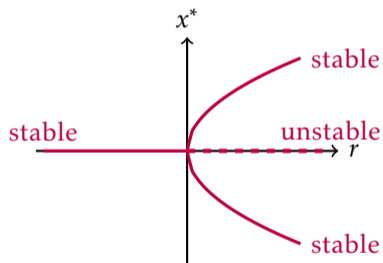


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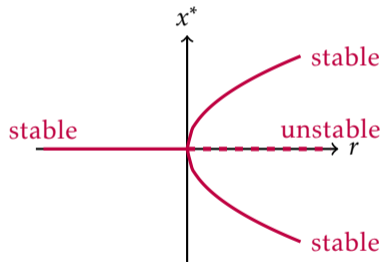
A stable equilibrium becomes unstable and two new stable equilibria branch out of it.

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(this is also why these types of bifurcations are called *pitchfork bifurcations*).





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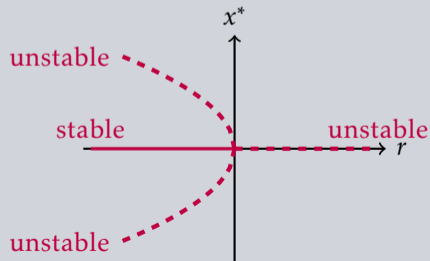
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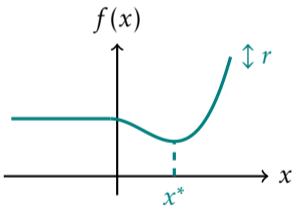
I leave you as an exercise to study this case (it's the same as before with a different sign). Draw a streamplot of the system for  $r < 0$ ,  $r = 0$  and  $r > 0$  and verify that the bifurcation diagram looks like this:



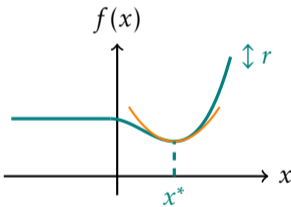


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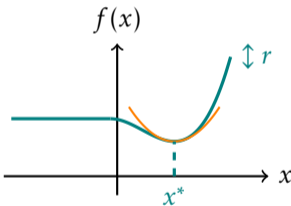


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Locally around  $x^*$ ,  $f(x)$  “looks like” a parabola of the form  $r + x^2$ , so we know that as  $r$  changes the system will exhibit a *saddle-point bifurcation* around  $x^*$ .

That's all!

Questions?

UC San Diego